

MHD in non-inductive tokamak plasmas: simulations and comparison to experiments on Tore Supra

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Aknowlegments to

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□ Non-inductive tokamak scenario: continuous energy production

- Unlimited plasma discharge
- Plasma current driven by
 - Self-generated **bootstrap current** (~15% on Tore Supra, more than 50% needed for ITER)
 - **Auxilliary systems** (~85% provided by Lower Hybrid (LH) waves on Tore Supra)

□ Non-inductive discharges tend to have hollow current profiles

- On Tore Supra: hollow current profile driven by **LH deposition**
- On ITER: hollow current profile driven by **off-axis bootstrap**

□ Magneto-Hydro-Dynamic stability of hollow current profiles

- At **moderate** $\beta = \mu_0 p / B^2$: (Double-)Tearing instabilities
- At **high** β : pressure driven modes (Neoclassical Tearing, Interchange, Ballooning, kink)

- ❑ **Impact of MHD instabilities on Tore Supra non-inductive pulses**
 - **Loss of wave-driven fast electrons**
 - Less current drive
 - Risks for plasma facing components (water leaks...)
 - **Degradation of energy confinement**

- ❑ This has motivated **studies aiming at**
 - Understanding **what determines MHD-stable domain**
 - Understanding **how MHD modes impact energy confinement**
 - *Understanding **how Tore Supra case extrapolates to higher- β***

- ❑ **Tools**
 - **MHD diagnostics** (δT_e , δB)
 - **Numerical codes for linear and non-linear studies**

Outline

- Numerical tools
- Impact of the Double-Tearing Mode
- Fully developed non-linear regimes
- Issue on $n=1$ mode stability
- Conclusion and perspectives

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- **Tool for non-linear MHD:** XTOR [Lütjens '08] (full MHD, toroidal)

Toroidal rotation

$$\rho (\partial_t \mathbf{v} + \mathbf{v} \cdot \nabla \mathbf{v}) = \mathbf{J} \times \mathbf{B} - \nabla p + \mu_\varphi (\mathbf{V}_{\text{src}} - V_\varphi) + \nu \nabla^2 \mathbf{v}$$

$$\partial_t p + \mathbf{v} \cdot \nabla p = -\Gamma p \nabla \cdot \mathbf{v} + \nabla \cdot \chi_\perp \nabla p + \dots$$

$$\dots + \mathbf{B} \cdot \nabla [\chi_{//} (\mathbf{B} \cdot \nabla p) / B^2] + H$$

$$\partial_t \mathbf{B} = -\nabla \times \mathbf{E} \quad \text{Parallel transport}$$

$$\mathbf{E} + \mathbf{V} \times \mathbf{B} = \eta (\mathbf{J} - \mathbf{J}_{\text{NI}})$$

$$\mathbf{J}_{\text{NI}} = \mathbf{J}_{\text{cd}} + \mathbf{J}_{\text{boot}} \quad \text{Non-inductive sources}$$

$$\mathbf{J}_{\text{cd}} = (\mathbf{J} - \mathbf{J}_{\text{bs}})_{t=0}$$

- **Standard MHD model adapted for Tore Supra experiments**

- **Non-inductive current sources**
- **Anisotropic Heat transport**
- **Toroidal rotation source (as generated by ripple)**

□ Re-Scaling of dimensionless coefficients

- Lundquist number ($S = \tau_R / \tau_A$) $\sim 10^9$ in experiments
- Hardly tractable in numerical simulations for resistive MHD modes
 - Fine radial discretisation ($\Delta r \sim S^{-2/5}$)
 - Long time evolution ($T \sim S^{3/5}$)
- Using lower S implies rescaling other parameters
 - **Option MHD / pressure dynamics**
 - ✓ Correct pressure dynamics during mode growth
 - ✓ $\tau_{\text{MHD}} / \tau_E = (\tau_{\text{MHD}} / \tau_E)^{\text{exp}}$: $S^{3/5} \chi_{\perp} = (S^{3/5} \chi_{\perp})^{\text{exp}}$
 - **Option current / pressure dynamics**
 - ✓ Correct pressure dynamics during equilibrium evolution
 - ✓ $\tau_R / \tau_E = (\tau_R / \tau_E)^{\text{exp}}$: $S \chi_{\perp} = (S \chi_{\perp})^{\text{exp}}$
- Typically, we use $S = 5 \times 10^6 - 5 \times 10^7$

□ Other parameters

- Anisotropic diffusivity: $\chi_{\parallel} / \chi_{\perp} = 10^8$
- Prandtl number : $S_v = 1$

Outline

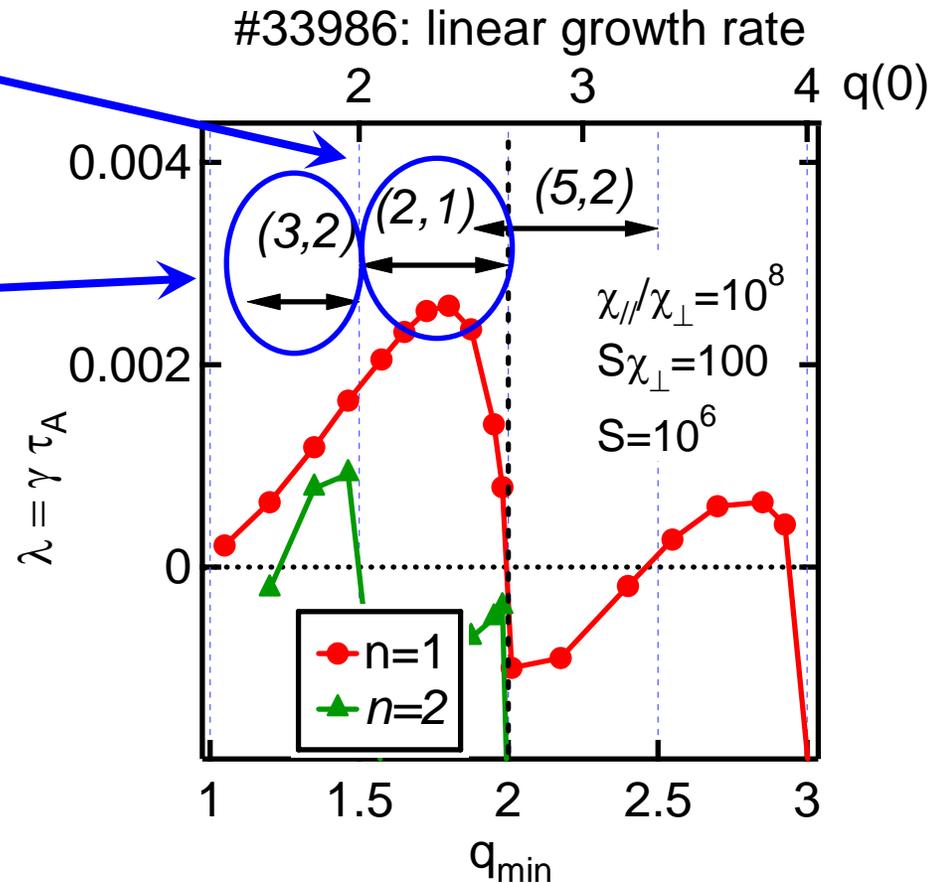
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Basic linear MHD stability features

- q-scan by Rescaling of total current
- Tearing unstable when q_{\min} below 2

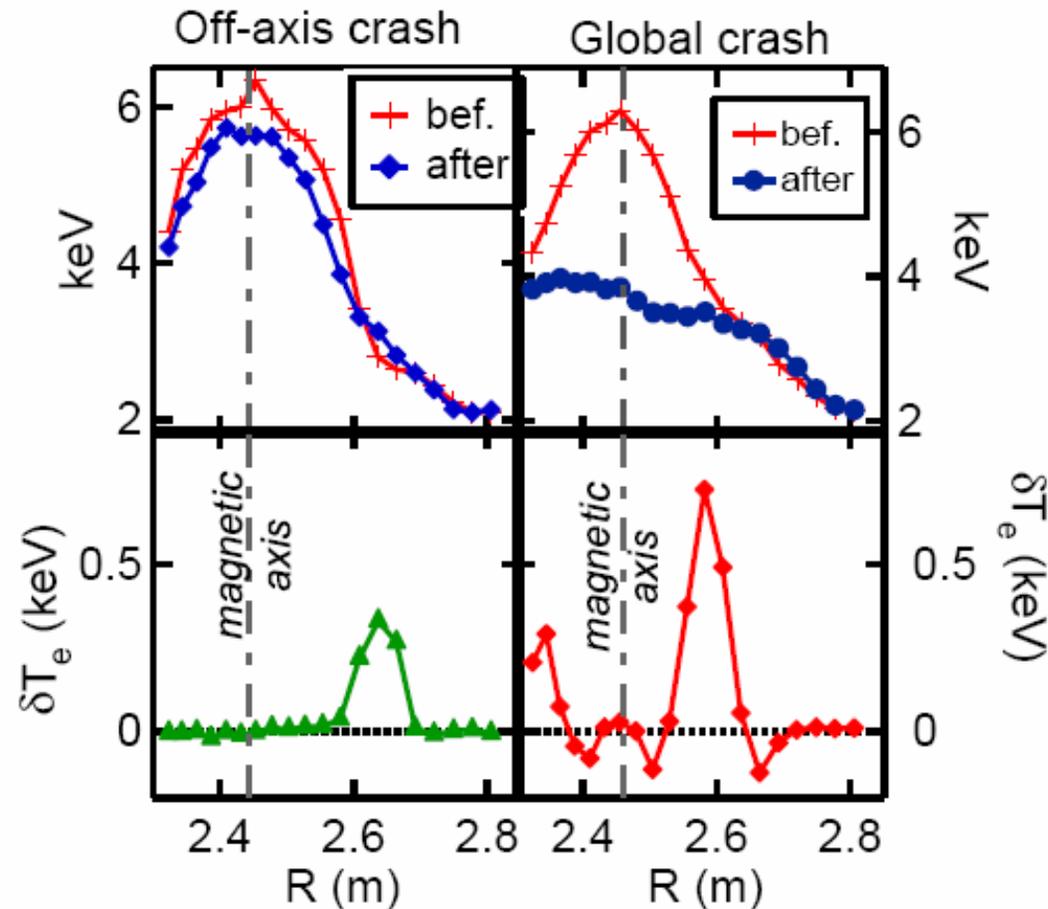
DTM on $q=2/1$

DTM on $q=3/2$



- Impact of the Double-Tearing Mode (DTM) on confinement
 - Observation of different situations
 - Off-axis crashes, complete crash (q=2 DTM)

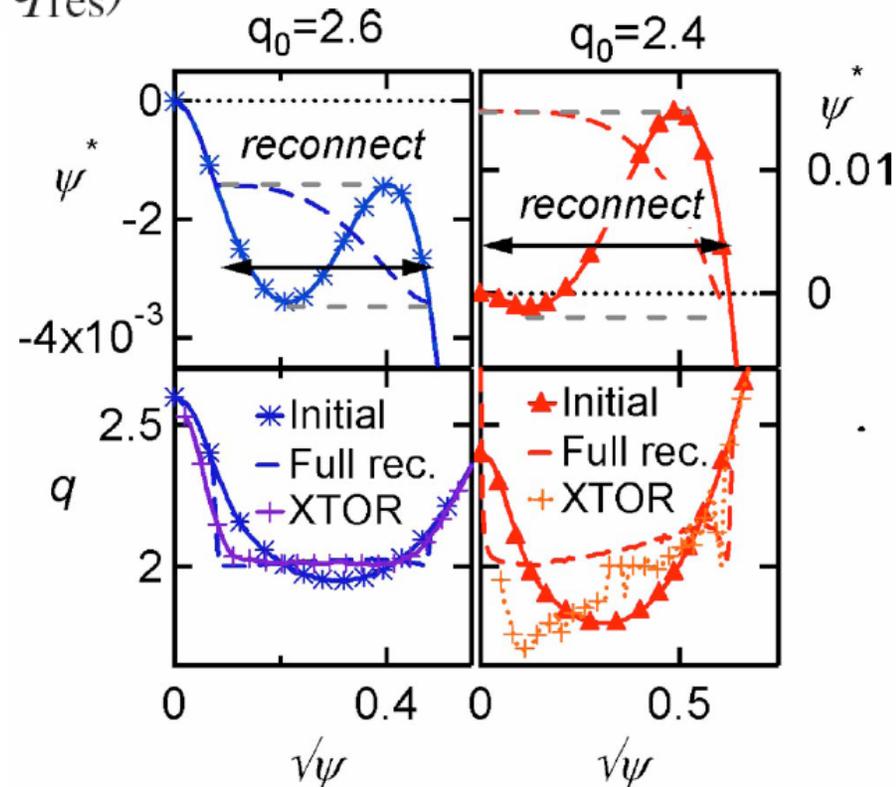
What determines the impact of the DTM ?



□ Unstable DTM experiences full reconnection

- Full reconnection model [Kadomtsev '75] extended to non-monotonic q-profiles [Carreras '79]
 - Helical flux reconnected in multiple-valued region

$$\psi^* = \int d\psi_N (1 - q/q_{res})$$



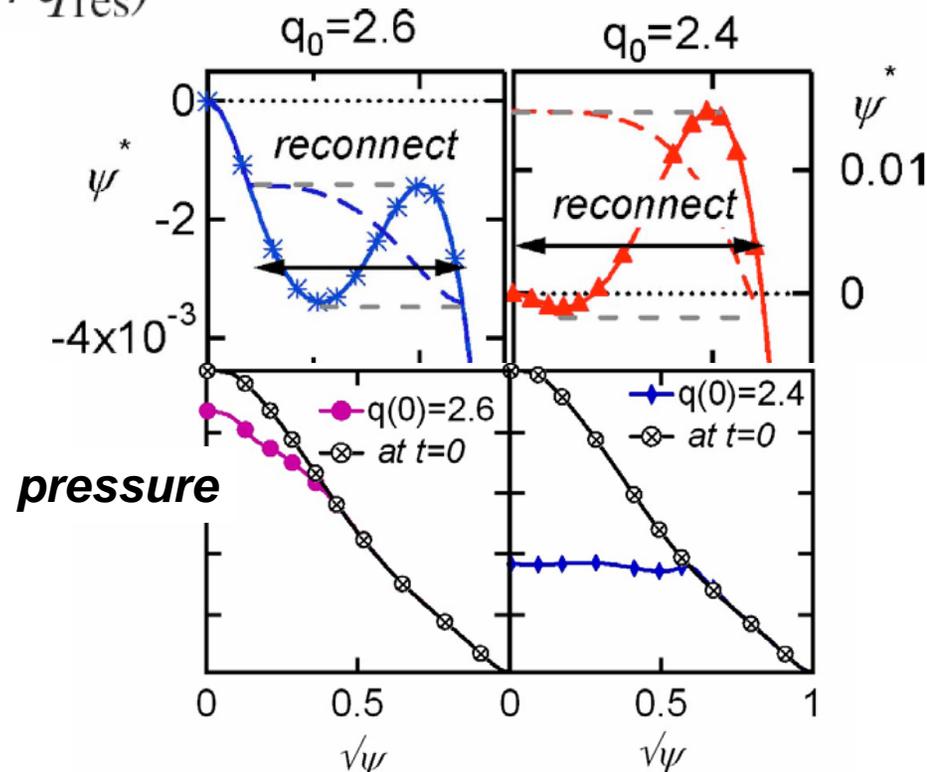
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$$\psi^* = \int d\psi_N (1 - q/q_{res})$$

- Core confinement
 - Small impact if off-axis
 - Else, global crash

Off-axis and global reconnections at $q=2$ consistent with full reconnection model and XTOR simulations



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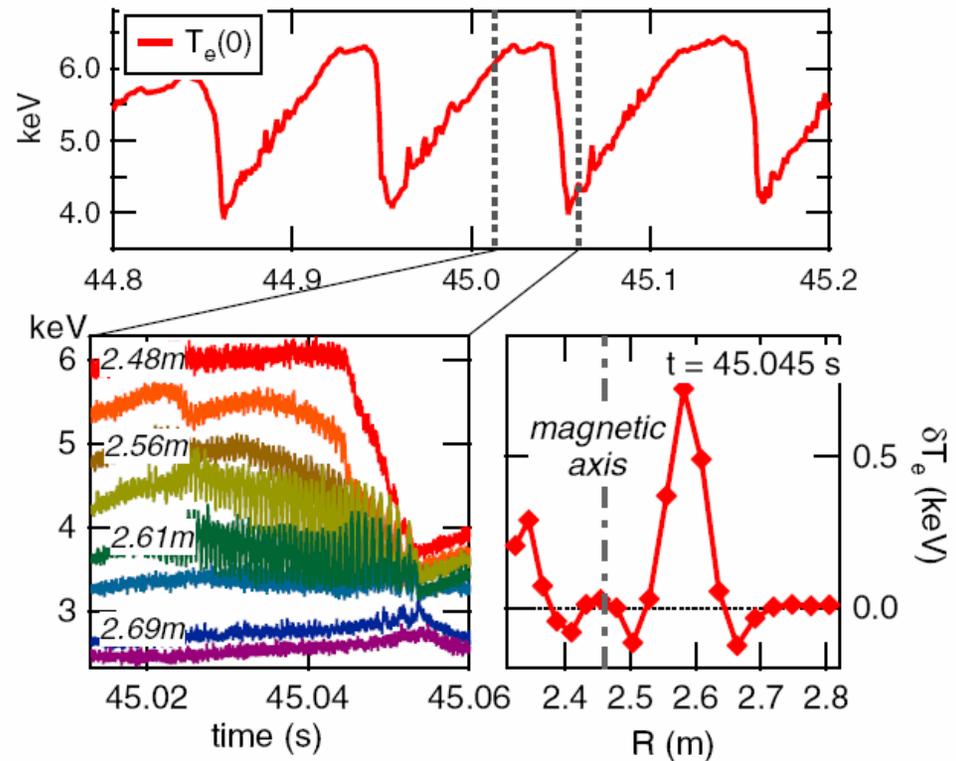
□ What is the dynamics of MHD-unstable plasma ?

- DTM can produce periodic relaxations

- The $q=2$ sawtooth regime [Chang '96]

- *Example from Tore Supra*

*$q=2$ sawtooth regime
with global crashes
(consistent with full
reconnection)*

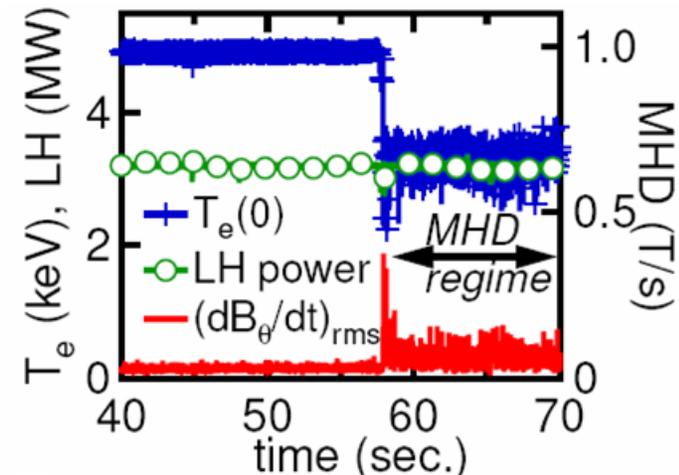


□ What is the dynamics of MHD-unstable plasma ?

- **DTM can produce periodic relaxations**
 - The $q=2$ sawtooth regime [Chang '96]

- **In other conditions, saturated regime with MHD activity**
 - Confinement degradation
 - Losses of fast electrons (LH-driven)

Transition to a regime with saturated MHD



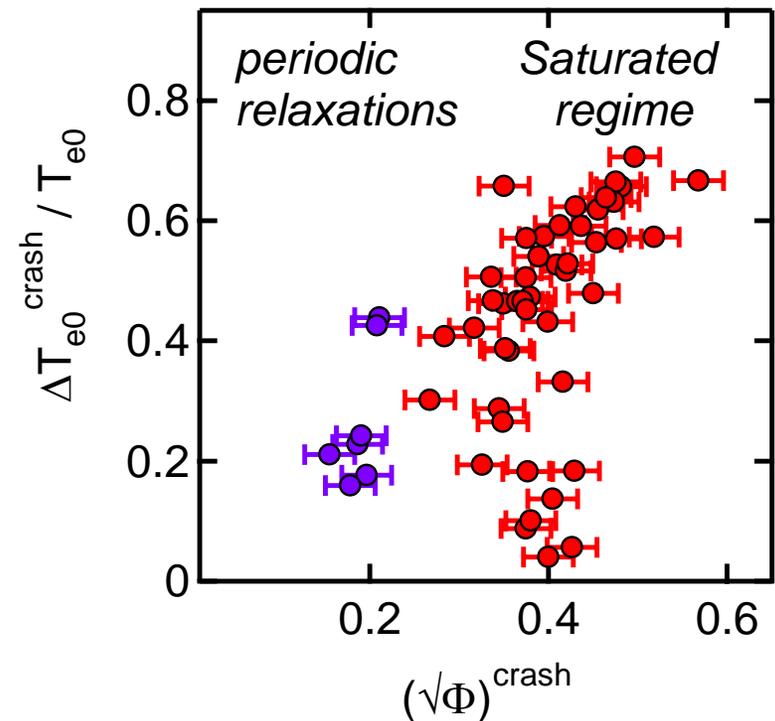
- Discrimination between periodic relaxations and saturation
 - Experimental database for Tore Supra shows
 - Periodic regime if crash radius small enough

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*Non-linear regimes
and position of
inversion radius after
MHD mode*

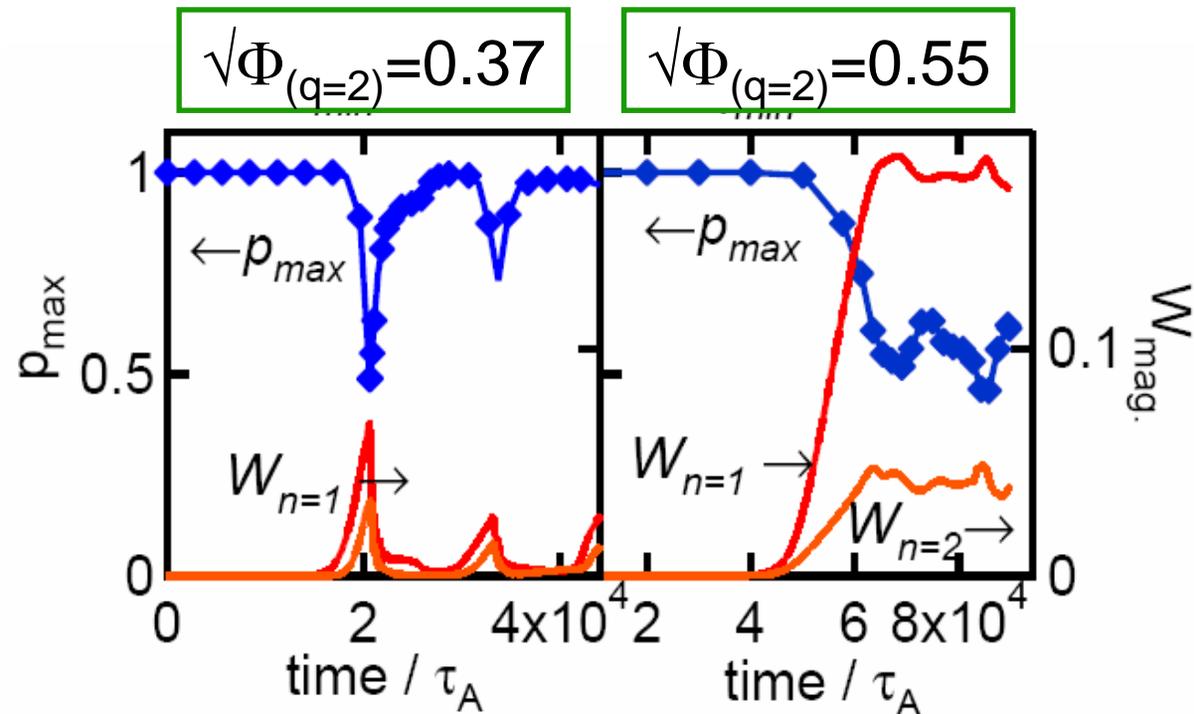


□ Non-linear MHD simulations

- Here $\tau_R/\tau_E = (\tau_R/\tau_E)^{\text{exp}}$ to have consistent pressure / equilibrium dynamics
- **Total current rescaled to increase $q=2$ radius**
 - *At the same time, goes from Double to Single tearing*

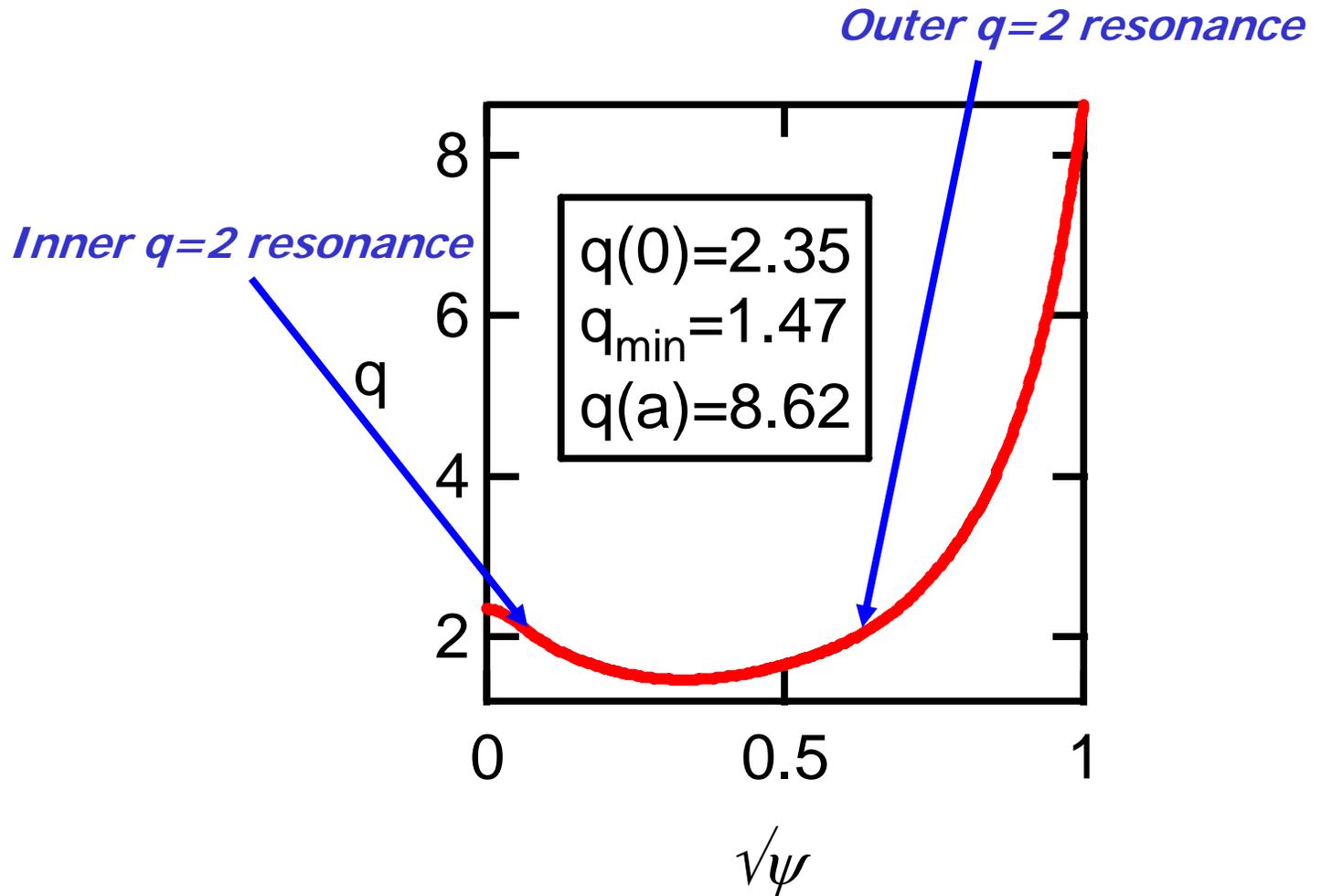
Non-linear regimes for different $q=2$ positions

- *periodic relaxations for DTM*
- *saturated state when $q=2$ more outside and STM*



□ Fully non-inductive discharges ($q_{\min} \sim 3/2$):

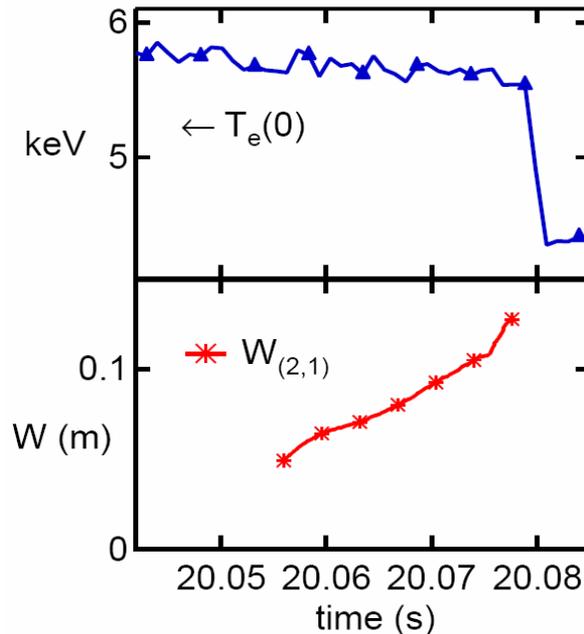
- Magnetic equilibrium



□ Fully non-inductive discharges ($q_{\min} \sim 3/2$): two regimes

- Successfull long discharges
 - No sign of $n=1$ activity
- But in similar conditions
 - $n=1$ mode can grow and saturate at large size

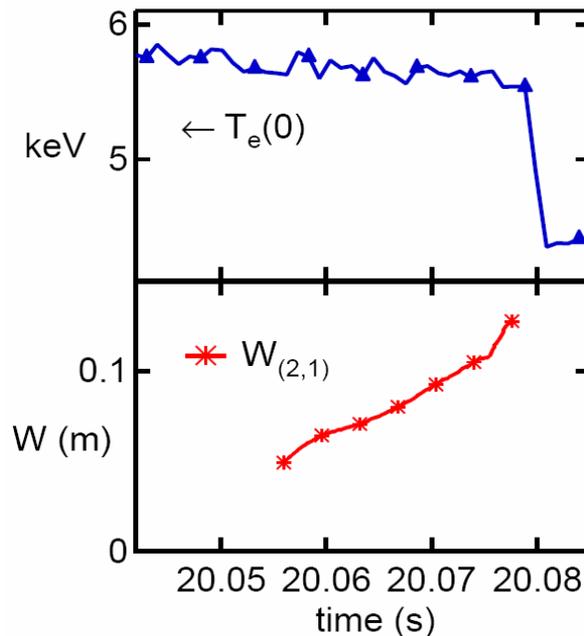
Example of $n=1$ becoming unstable



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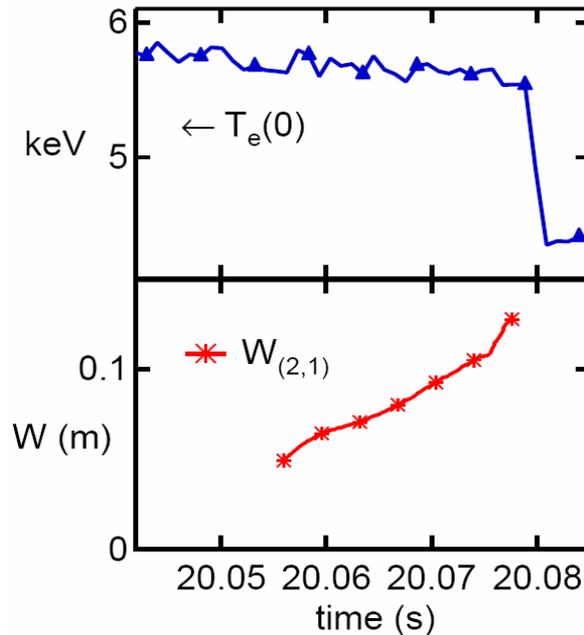
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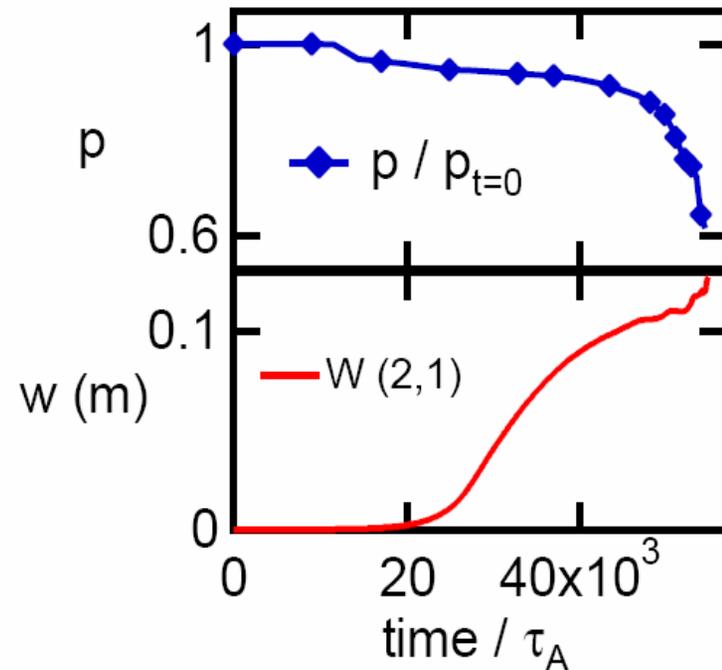
Can non-linear MHD model recover these two regimes ?

- Long fully non-inductive discharges: wrong $n=1$ prediction
 - *Non-linear MHD simulations including transport in agreement with global crash: $n=1$ mode found unstable*
 - *Successfull discharges (no $n=1$ mode) cannot be recovered*

Example of $n=1$ becoming unstable



XTOR $S=5 \times 10^6$



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□ Condition for recovering n=1 stability in standard MHD

- Curvature term at high S [*Glasser '75 (GGJ)*]
- No transport

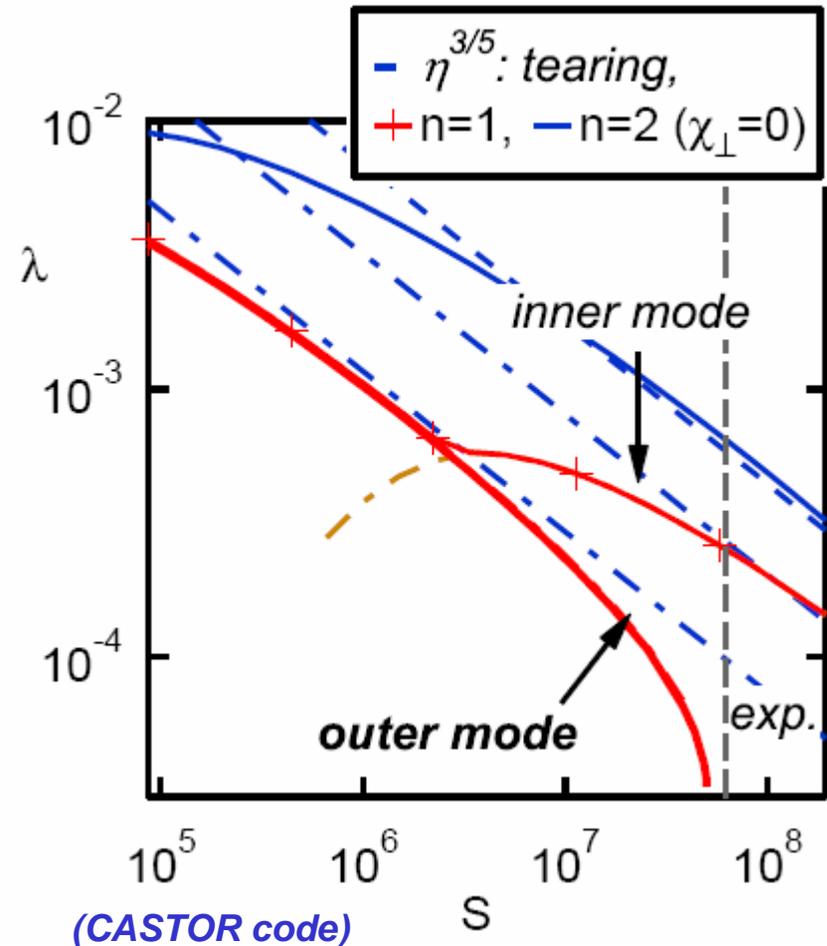
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2 tearing branches (2 resonances)

Outer mode stable at $S=S^{exp}$



□ Condition for recovering n=1 stability in standard MHD

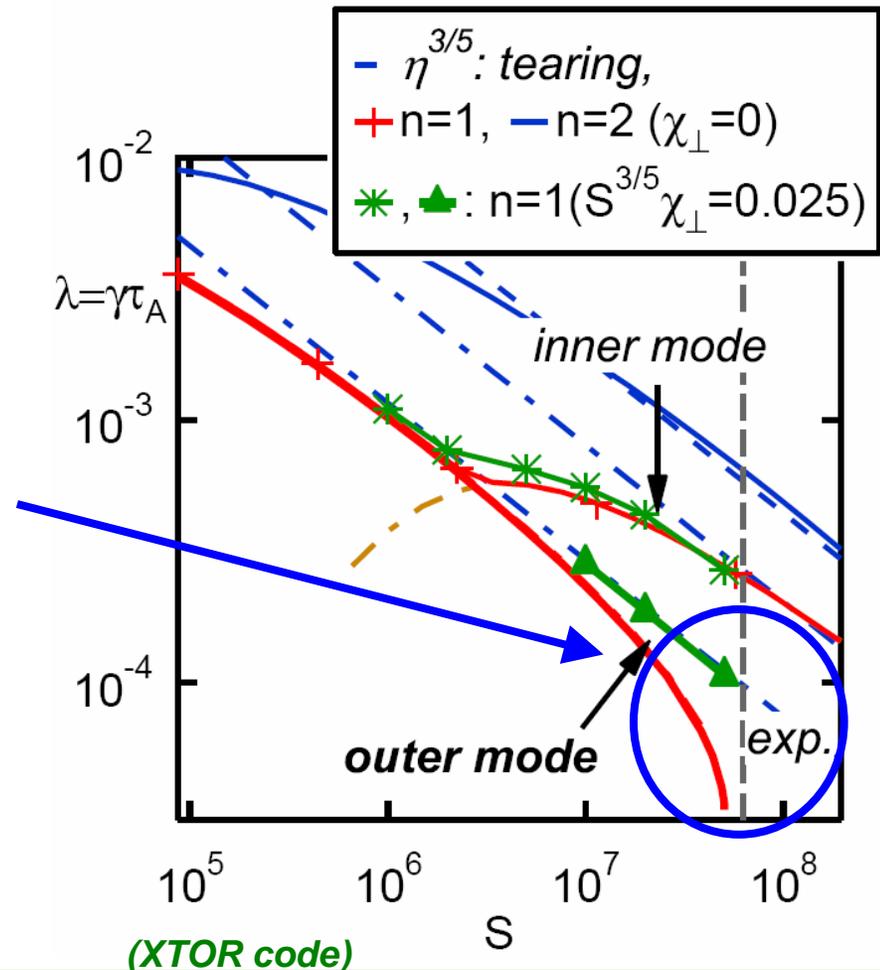
- Curvature term at high S [*Glasser '75 (GGJ)*]
- No transport

• **But transport is expected to play a role (low growth rate)**

n=1 then unstable at $S=S^{exp}$

• *consistent with theory expectations [Lütjens '01]*

Additional physics should be considered



□ Candidate mechanisms for n=1 stabilisation

- Toroidal rotation (ripple-driven in Tore Supra)
- Small island physics (outside standard MHD description)

□ Rotation effect

- Rotation in Tore Supra consistent with theory prediction for ripple drive [*Trier '08*]: Mach number is low ($M \sim 0.05$)

Toroidal rotation

$$\rho (\partial_t \mathbf{v} + \mathbf{v} \cdot \nabla \mathbf{v}) = \mathbf{J} \times \mathbf{B} - \nabla p + \mu_\varphi (\mathbf{V}_{\text{src}} - V_\varphi) + \nu \nabla^2 \mathbf{v}$$

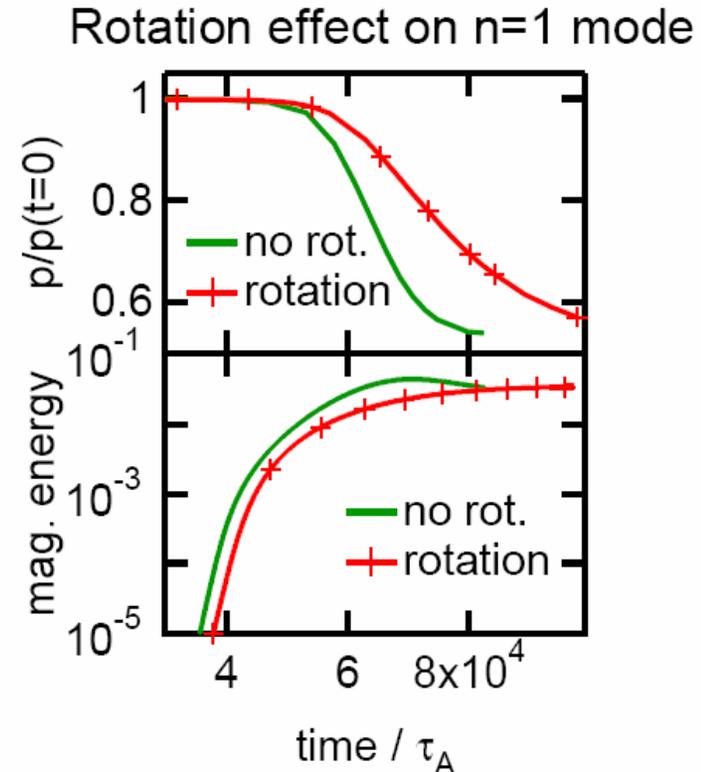
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- Toroidal rotation (ripple-driven in Tore Supra)
- Small island physics (outside standard MHD description)

□ Rotation effect

*No effect on (2,1)
tearing saturation*

*The pressure crash still
occurs*



□ Diamagnetic rotation: possible candidate

- Not negligible a priori: $\omega_e^* / \gamma \sim 10$
- Could potentially
 - Limit island growth in the non-linear regime [*Scott '87*]
 - Create multiple saturated states [*Ottaviani '04*]
- **Investigations**
 - 4-field reduced MHD, cylindrical geometry [*Hazeltine '87*]
 - Preliminary study with XTOR

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□ 4-field Reduced MHD [Hazeltine '87]

- Poloidal flux (ψ), electrostatic potential (ϕ), electron pressure (p), parallel ion velocity (v)
- Isothermal model

$$\begin{aligned} \partial_t \psi &= \eta \nabla_{\perp}^2 \psi - \nabla_{\parallel} (\phi + \delta p) + e \\ \partial_t W &= v \nabla_{\perp}^2 W + [W, \phi + \delta \tau p] - \nabla_{\parallel} J + \delta \tau [\nabla_{\perp} p; \nabla_{\perp} \phi] \\ \partial_t p &= \kappa_{\perp} \nabla_{\perp}^2 p + [p, \phi] - \beta \nabla_{\parallel} (v - 2\delta J) + S_p \\ \partial_t v &= v_{\parallel} \nabla_{\perp} v + [v, \phi] - \frac{1 + \tau}{2} \nabla_{\parallel} p + \delta \tau \beta \frac{1 + \tau}{2} [p, v] \end{aligned}$$

$$W = -\nabla_{\perp}^2 \phi \quad J = -\nabla_{\perp}^2 \psi$$

$$\nabla_{\parallel} \equiv \varepsilon \partial_{\phi} - [\psi, \cdot]$$

$$\omega_e^* = -(m/x) \delta(dp/dx) \text{ with } x = r/a$$

□ 4-field Reduced MHD [Hazeltine '87]

- Non-linear simulations with $n=0,1$
- Experimental resistivity $\eta=5 \times 10^{-9}$, $\nu=v_{||}=\eta$
- $\omega_e^*=4 \times 10^{-3}$ (higher than experimental $\omega_e^*=10^{-3}$)

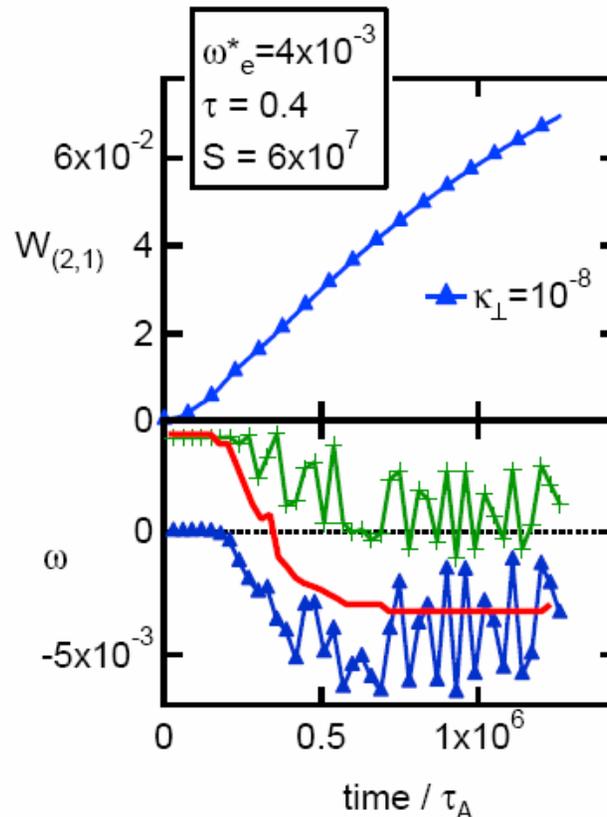
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Two non-linear regimes

- *higher κ_{\perp} :*
 - *large island*
 - *$\omega \sim \omega_{ExB}$*
 - *high ω_{ExB} generation*



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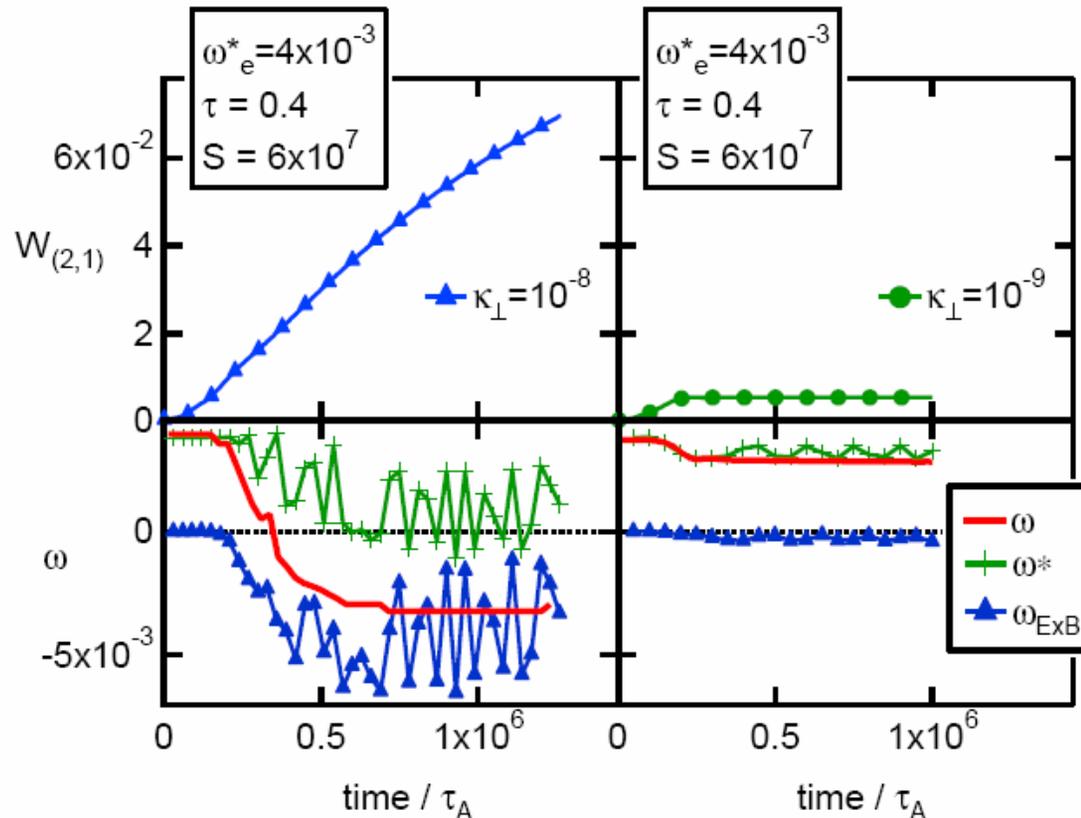
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Two non-linear regimes

- **higher κ_{\perp} :**
 - large island
 - $\omega \sim \omega_{ExB}$
 - high ω_{ExB} generation
- **lower κ_{\perp} :**
 - small saturated island,
 - $\omega \sim \omega^*$
 - low ω_{ExB} generation



□ 4-field Reduced MHD [Hazeltine '87]

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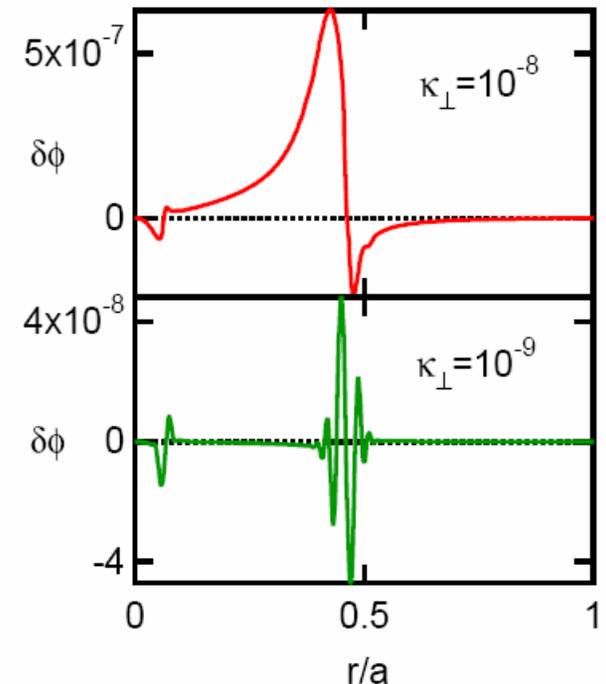


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Two non-linear regimes

- **higher κ_{\perp} :**
 - **large island**
 - $\omega \sim \omega_{ExB}$
 - **high ω_{ExB} generation**
- **lower κ_{\perp} :**
 - **small saturated island,**
 - $\omega \sim \omega^*$
 - **low ω_{ExB} generation**

**Distinct
radial
structures**



□ 4-field Reduced MHD [Hazeltine '87]

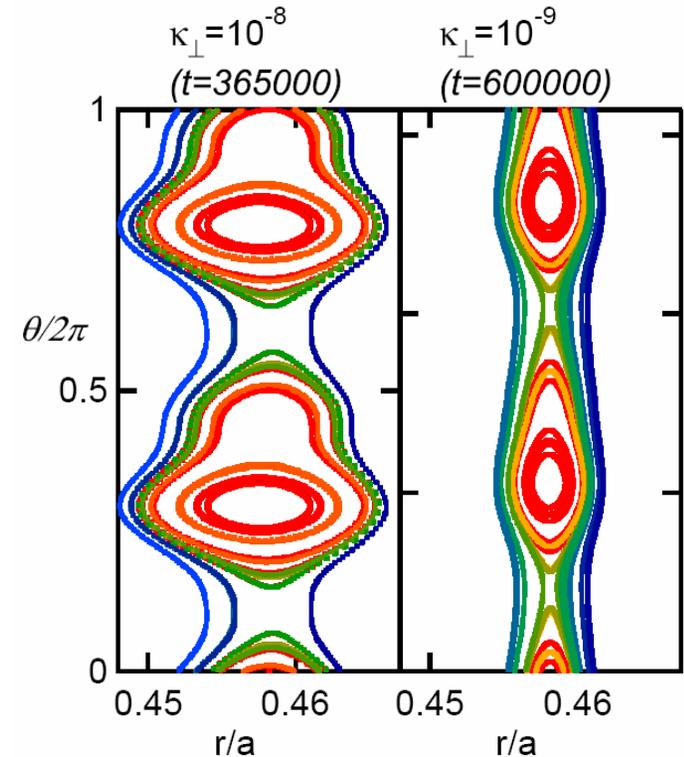
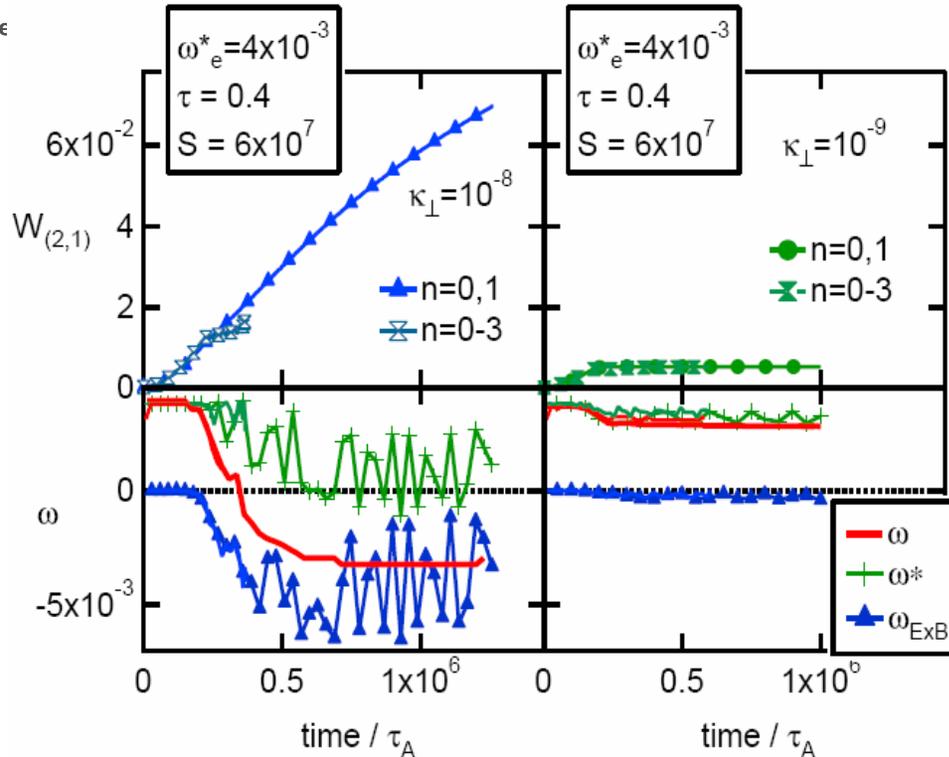
- Non-linear simulations (n=0,...3):

- same results
- Strong island distorsion
- Computation of the large island case crashes early

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□ 4-field Reduced MHD [Hazeltine '87]

- **Two different non-linear regimes**

- Controlled by κ_{\perp} and ω_e^*
- At low κ_{\perp} , pressure perturbation dominated by convection, otherwise dominated by transport
- Qualitatively satisfying for explaining experimental facts

- **Transition at $\kappa_{\perp} \in [10^{-9}, 4 \times 10^{-9}]$ for $\omega_e^* \sim 4 \times \omega_e^{* \text{exp}}$**

- κ_{\perp} consistent with collisional process
($\kappa_{\perp} = \eta(1 + \tau)/2 = 3.5 \times 10^{-9}$)
- But in experiment, $\kappa_{\perp} \sim 10^{-7}$: **large island regime**

- Finally, the question of transport remains central even when diamagnetic effects are considered

□ Diamagnetic rotation: toroidal effects

- Can toroidal geometry add significant stabilisation to n=1 mode ?
- Preliminary study with XTOR
- New version of the code: fully implicit time advance
- Electron diamagnetic effect in Ohm's law:

$$\mathbf{E} + \mathbf{V} \times \mathbf{B} = \eta (\mathbf{J} - \mathbf{J}_{\text{NI}}) - \frac{1}{1 + \tau} \frac{\nabla_{\parallel} p}{en}$$

□ Diamagnetic rotation: toroidal effects

- Linear study with transport

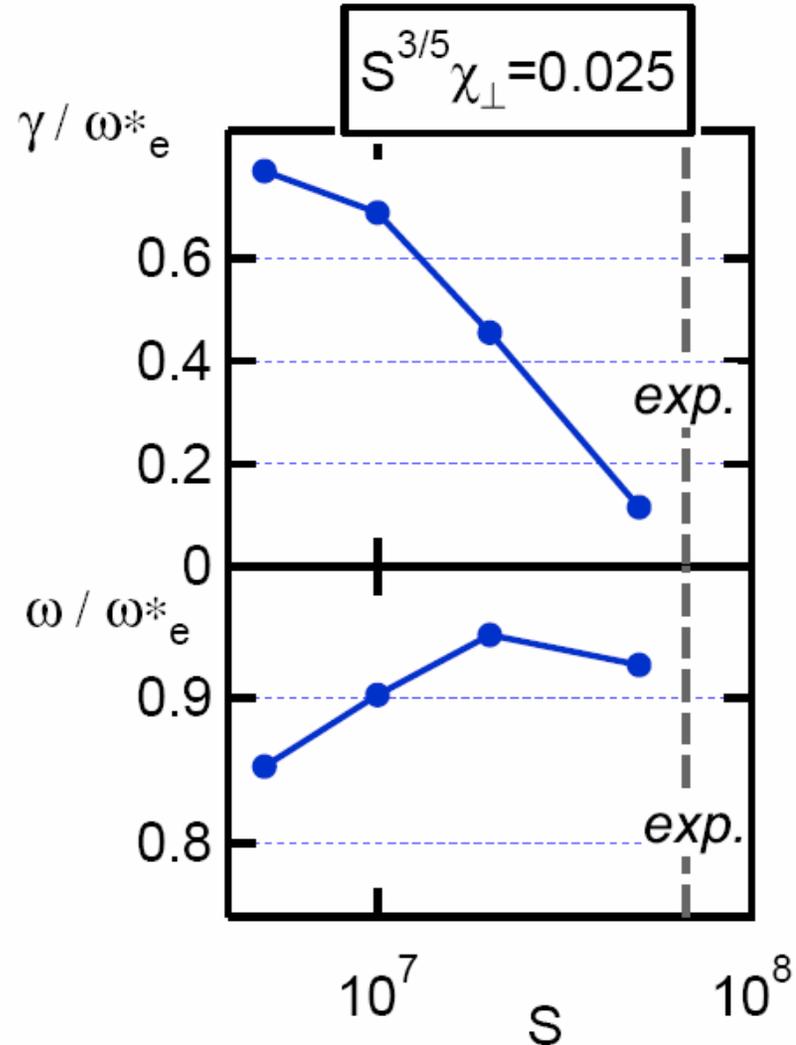
➤ $S^{3/5} \chi_{\perp} = (S^{3/5} \chi_{\perp})^{exp}$

➤ At experimental S

✓ $\gamma \ll \omega^*_e$

✓ $\omega \sim 0.95 \omega^*_e$

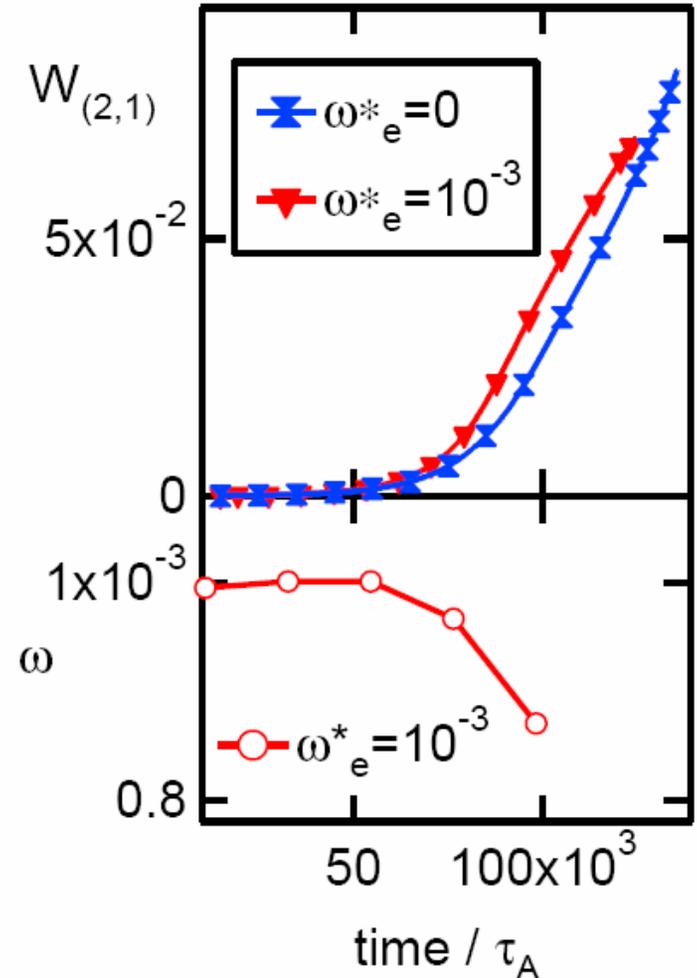
We expect significant effect of diamagnetic rotation in the full MHD model including transport



□ Diamagnetic rotation: toroidal effects

- Non-linear regime with $n=0,1$
 - $S=5 \times 10^7$, $\gamma/\omega_e^* \sim 0.16$

XTOR, $S=5 \times 10^7$, $S^{3/5} \chi_{\perp} = 0.025$



No significant difference compared to $\omega_e^=0$ case*

Consistent with reduced-MHD 4-field model:

- *small island regime not accessible in presence of realistic transport coefficients*

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- ❑ **Dynamics of MHD modes in non-inductive discharges**
 - **Double-Tearing: full reconnection**
 - **Saturation / Periodic relaxations: $q=2$ surface position**

- ❑ **Stability of $n=1$ mode in very long discharges**
 - **Curvature stabilization at high S without transport**
 - **But fails in more realistic conditions (finite transport)**
 - **Two-fluid effect allows small island regime**
 - **But at collisional transport coefficient level only**

- ❑ **Future work:**
 - more comprehensive MHD model needed ?
 - Other rationale for transport coefficient implementation ?